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# Calculation of the elastic properties of semiconductors

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Abstract. For 30 diamond- and zincblende-structure semiconductors, the bond length d, bond polarity  $\alpha_p$ , bulk modulus B, elastic shear constants  $(c_{11} - c_{12})/2$  and  $c_{44}$ , bond-stretching force constant  $\alpha$ , bond-bending force constant  $\beta$ , internal displacement parameter  $\zeta$ , effective atomic charge  $Z^*$ , transfer parameter  $\beta^*$ , transverse charge  $e_T^*$ , and piezoelectric charge  $e_P^*$  are calculated from bond orbital calculations based on the tight-binding method. The results are compared with previous theoretical calculations and experiments.

### **1. Introduction**

Theoretical studies for the elastic properties of semiconductors have been made by many researchers, using, for example, a bond-orbital model (BOM) [1,2], band-structure calculations and the 'special-points' method [3], a first-principles linear-combination of atomic orbitals (LCAO) method [4], pseudopotential methods with the local-density approximation (LDA) [5-7], the linear muffin-tin orbitals (LMTO) method [8,9], the universal tight-binding parameters (UTBP) method [10], and the extended Hückel tight-binding (XHTB) method [11].

The LDA has proven to be an effective and useful means for studying both structural and electronic properties in a few materials [5–7]. However, the inclusion of a total-energy calculation in the problem necessarily makes the LDA theory more complicated than the BOM theory; systematic study of the properties in many materials would need a powerful computer, and it will take much time; therefore, the cost will be expensive.

In order to obtain the structural trends, the simple BOM is used in this paper to study the elastic properties of semiconductors systematically. We calculate the bond length d, and bond polarity  $\alpha_p$ , bulk modulus B, elastic shear constants  $(c_{11} - c_{12})/2$  and  $c_{44}$ , bondstretching force constant  $\alpha$ , bond-bending force constant  $\beta$ , internal displacement parameter  $\zeta$ , effective atomic charge  $Z^*$ , transfer parameter  $\beta^*$ , transverse charge  $e_T^*$ , the piezoelectric charge  $e_p^*$ , and the ratios  $c_{12}/c_{11}$ ,  $[(c_{11}-c_{12})/2]/B$  and  $\beta/\alpha$  for 30 diamond- and zincblendestructure semiconductors and compare the values with others from theoretical calculations and experiments.

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## 2. Theoretical formalism

#### 2.1. Bond length and polarity

Tight-binding theory has been able to obtain approximate but meaningful predictions of the bonding properties of solids. In order to improve the description of the bond energy, Baranowski [2] made a very simple modification of the overlap interaction on the basis of the idea and method proposed by Harrison [1] and gave a formula for the bond length of semiconductors. For all tetrahedral compounds, the bond length d can be obtained from [2]

$$d = \frac{(-2\eta_{a}\hbar^{2}/m)^{1/2}}{(k^{2}\bar{\varepsilon}_{b}^{2} - 4V_{3}^{2})^{1/4}}.$$
(1)

where for sp<sup>3</sup> bonds  $\eta_{\sigma} = \frac{1}{4}\eta_{ss\sigma} - (2\sqrt{3}/4)\eta_{sp\sigma} - \frac{3}{4}\eta_{pp\sigma}$ , in which  $\eta_{ss\sigma} = -1.4$ ,  $\eta_{sp\sigma} = 1.84$  and  $\eta_{pp\sigma} = 3.24$  are dimensionless Harrison [1] universal parameters and  $\hbar^2/m = 7.62$  eV Å<sup>2</sup>. The effective parameter k will be given by the following average:

$$k = (k_i k_i)^{1/2}$$
(2)

where  $k_i$  and  $k_j$  are connected with rows *i* and *j*, respectively, of the periodic table. The cation-anion average hybrid energy  $\bar{e}_h$  is the weighted average given by

$$\bar{\varepsilon}_{\rm h} = \frac{1}{8} (n_{\rm c} \varepsilon_{\rm h}^{\rm c} + n_{\rm a} \varepsilon_{\rm h}^{\rm a}) \tag{3}$$

where  $n_c$  and  $n_a$  are the numbers of electrons associated with cations and anions, respectively, which participate in the bonds.  $\varepsilon_h^c$  and  $\varepsilon_h^a$  are the average values of the cation and anion hybrid energies, respectively:

$$\varepsilon_{h}^{c} = \frac{1}{4} (\varepsilon_{s}^{c} + 3\varepsilon_{p}^{c})$$

$$\varepsilon_{h}^{a} = \frac{1}{4} (\varepsilon_{s}^{a} + 3\varepsilon_{p}^{a})$$
(4)

in which  $\varepsilon_s^c$ ,  $\varepsilon_p^c$ ,  $\varepsilon_s^a$  and  $\varepsilon_p^a$  are the free-atom energies for s and p states for the cations and anions, respectively [12].  $V_3$  is the hybrid polar energy, which can be approximated in the following way:

$$V_3 = \frac{1}{2} (\varepsilon_h^c - \varepsilon_h^a). \tag{5}$$

The bond polarity of the compound can be obtained from

$$\alpha_{\rm p} = \frac{V_3}{(V_2^2 + V_3^2)^{1/2}} \tag{6}$$

where  $V_2$  is the hybrid covalent energy, which can be approximated in the following way:

$$V_2 = -\eta_\sigma \frac{\hbar^2}{md^2}.$$
(7)

## 2.2. Bulk modulus and elastic shear constant

According to the BOM proposed by Baranowski [2], the bulk modulus is given by

$$B = \frac{2\sqrt{3}}{3} \left( V_2 \alpha_{\rm c}^3 + \frac{7.8}{d^2} \right) \frac{1}{d^3}$$
(8)

where the hybrid covalency  $\alpha_c$  is defined by  $(1 - \alpha_p^2)^{1/2}$ .

The elastic shear constant  $(c_{11} - c_{12})/2$  is given by

$$\frac{(c_{11} - c_{12})}{2} = \frac{\sqrt{3}}{4} (V_2 \alpha_c^3 (1 + \lambda) - \frac{3}{4} |V_{pp\pi}|) \frac{1}{d^3}$$
(9)

where  $\lambda$  is a dimensionless parameter, which is defined as follows:

$$\lambda = \frac{\sqrt{3}V_{\rm sp\sigma} - 3V_{\rm pp\sigma}}{V_{\rm ss\sigma} - 2\sqrt{3}V_{\rm sp\sigma} - 3V_{\rm pp\sigma}}$$
(10)

where  $V_{ss\sigma}$ ,  $V_{sp\sigma}$ ,  $V_{pp\sigma}$  and  $V_{pp\pi}$  in equation (9) are the matrix elements [1]. Combining equations (8), (9) and  $B = \frac{1}{3}(c_{11} + 2c_{12})$ , one can obtain the elastic constants

$$c_{11} = \frac{\sqrt{3}}{3} \left( V_2 \alpha_{\rm c}^3 (3+\lambda) - \frac{3}{4} |V_{\rm pp\pi}| + 2\frac{7.8}{d^2} \right) \frac{1}{d^3}$$
(11)

and

$$c_{12} = \frac{\sqrt{3}}{6} \left( V_2 \alpha_c^3 (3 - \lambda) + \frac{3}{4} |V_{pp\pi}| + 4\frac{7.8}{d^2} \right) \frac{1}{d^3}.$$
 (12)

Within the framework of the valence force field model [1], the elastic stiffness constant  $c_{44}$  and Kleinman's internal displacement parameter  $\zeta$  of diamond- and zincblende-structure semiconductors are given by

$$c_{44} = 3(c_{11} + 2c_{12})(c_{11} - c_{12})/(7c_{11} + 2c_{12})$$
(13)

$$\zeta = (c_{11} + 8c_{12})/(7c_{11} + 2c_{12}). \tag{14}$$

The macroscopic elastic constants B and  $(c_{11} - c_{12})/2$  are related to force constants  $\alpha$  and  $\beta$  by the equations [13]

$$3B = \frac{\sqrt{3}}{4d}(3\alpha + \beta) - 0.355SC_0 \tag{15}$$

and

$$\frac{(c_{11} - c_{12})}{2} = \frac{\sqrt{3}}{2d}\beta - 0.053SC_0 \tag{16}$$

where  $SC_0$  is the Coulomb contribution (rather small for all III–V zincblende-structure semiconductors). If we neglect the Coulomb contribution, the bond-stretching force constant  $\alpha$  and bond-bending force constant  $\beta$  can be obtained from

$$\beta = \frac{2d}{\sqrt{3}} \frac{(c_{11} - c_{12})}{2} \tag{17}$$

$$\alpha = \frac{4d}{\sqrt{3}}B - \frac{1}{3}\beta. \tag{18}$$

		d (Å)			d (Å)
Material	BOM	Experiment	Material	BOM	Experiment
c	1.615	1.54	InAs	2.618	2.61
Si	2.357	2.35	InSb	2.820	2.81
Ge	2.448	2.44	BeS	1.900	2.10
Sn	2.802	2.81	BeSe	1.988	2.20
SiC	1.953	1.88	BeTe	2.190	2.40
BN	1.573	1.57	ZnS	2.302	2.34
BP	1.957	1.97	ZnSe	2.413	2.45
BAs	2.016	2.07	ZnTe	2.658	2.64
AlP	2.343	2.36	CdS	2.478	2.53
AlAs	2.417	2.43	CdTe	2.865	2.80
AISb	2.610	2.66	CuF	1.534	1.84
GaP	2.370	2.36	CuCi	2.122	2.34
GaAs	2.443	2.45	CuBr	2.268	2.49
GaSb	2.635	2.65	CuI	2.576	2.62
InP	2.538	2.54	AgI	2.799	2.80

Table 1. Predicted bond length of semiconductors. The experimental bond lengths are taken from [14].

### 2.3. Transverse and piezoelectric charges

An effective charge  $Z^*$  for a cation can be obtained in tight-binding theory as the column number Z (it is less than four for a cation) for that element minus the sum over occupied states of the squared amplitudes of the orbitals of that atom. The bond orbital approximation gives an approximate value of [1]

$$Z^* = Z - 4 + 4\alpha_{\rm p}.\tag{19}$$

The transverse charge  $e_T^*$  is given by

$$e_{\rm T}^* = Z^* + \frac{4}{3}\beta^* \tag{20}$$

where  $\beta^*$  is the transfer parameter, which is defined by

$$\beta^* = (d/4)\partial Z^*/\partial d = 2\alpha_{\rm p}(1 - \alpha_{\rm p}^2). \tag{21}$$

Similarly, the piezoelectric charge  $e_p^*$  is given by

$$z_{\rm p}^* = Z^* - (4\beta^*/3)(1-\zeta)/\zeta.$$
 (22)

All effective charges for the anion are the negative charges of the corresponding charges for the cations given here.

## 3. Results and discussion

In this section we present the results of the BOM calculation for the bond length d, bond polarity  $\alpha_p$ , bulk modulus B, elastic shear constants  $(c_{11} - c_{12})/2$  and  $c_{44}$ , bondstretching force constant  $\alpha$ , bond-bending force constant  $\beta$ , internal displacement parameter  $\zeta$ , effective atomic charge  $Z^*$ , transfer parameter  $\beta^*$ , transverse charge  $e_T^*$  and piezoelectric charge  $e_p^*$  for 30 diamond- and zincblende-structure semiconductors and compare their values with other values from theoretical calculations and experiments.

(i) Predictions of the bond length for semiconductors are given in table 1 together with the experimental values [14]. It is remarkable that the simple equation (1), which has been given by Baranowski [2], can reproduce the bond length with an accuracy of the order of a few per cent for most cases.

(ii) Results for the polarity  $\alpha_p$  obtained from BOM are listed in table 2, together with those for the UTBP, Hückel tight-binding (HTB), the XHTB methods [11], the cluster (Cl) method, and Brillouin-zone (BZ) integration of the LCAO Hamiltonian [15]. For comparison, the values of polarity  $\alpha_p (\equiv f_i^{1/2})$  given by Phillips (Ph) [11] and obtained by experiment (from the experiment-deduced ionicity scale (the definition is the same as Harrison's polarity) of Falter *et al* [16]), are also given in table 2. From a comparison of these values, we see that results obtained from BOM are in good agreement with the experiments.

Table 2. Comparison of polarities from various calculations. Results for polarity  $\alpha_p$  obtained from the BOM, the band-structure calculations based on the UTBP, HTB and XHTB methods, CI method, BZ integration of LCAO Hamiltonian. For comparison, the values of polarity  $\alpha_p$  obtained by Ph and in experiment are also given.

_					$\alpha_{p}$			
Material	CI	BZ	UTBP	HTB	XHTB	Ph	BOM	Experiment
SiC							0.16	0.10
AIP	0.42	0.44					0.34	0.33
AlAs	0.39	0.41					0.34	0.31
AISb	0.32	0.34				-	0.29	0.30
GaP	0.38	0.40	0.35	0.42	0.59	0.57	0.32	0.31
GaAs	0.35	0.37	0.33	0.37	0.53	0.56	0.32	0.30
GaSb	0.28	0.30	0.27	0.26	0.38	0.51	0.26	0.29
InP	0.46	0.48	0.43	0.46	0.66	0.65	0.40	0.32
InAs	0.42	0.44	0.40	0.42	0.61	0.60	0.40	0.30
InSb	0.35	0.37	0.35	0.33	0.49	0.57	0.34	0.29
ZnS	0.55	0.56	0.61	0.62	0.86	0.79	0.55	0.60
ZnSe	0.55	0.56	0.60	0.58	0.82	0.79	0.55	0.59
ZnTe	0.54	0.55	0.57	0.51	0.73	0.77	0.54	0.57
CdTe	0.58	0.59	0.63	0.54	0.77	0.82	0.61	0.58

(iii) The values of the transverse and piezoelectric charges obtained from equations (20) and (22) with a  $\beta$ -value obtained from equation (21), are listed together with the corresponding  $Z^*$  in table 3. The experimental values [16] are also listed for comparison. It is clear from table 3 that the effective atomic charges are in good agreement with the experiments in most cases; the transverse charges are about one electron less than in the experiments. When the experimental value of  $\alpha_p$  is used in equation (20), one cannot obtain the transverse charges, which are comparable with experiments. So the difference is not only from  $\alpha_p$ ; this suggests a significant error in equation (20). The piezoelectric charges are in good agreement with the experiments for III-V compound semiconductors.

(iv) The results for the covalency  $\alpha_c$ , the values of the bulk modulus *B* together with the elastic stiffness constants  $(c_{11} - c_{12})/2$  and  $c_{44}$ , bond-stretching force constant  $\alpha$ , bond-bending force constant  $\beta$  and Kleinman's internal displacement parameters  $\zeta$  are shown in table 4. The reliable values for  $\zeta$  are those obtained recently by Molinas-Mata and Cardona [17] by the fitting of the phonon dispersion curves along the [100] and [111] axes with planar force constants. The values that they obtained are  $0.564 \pm 0.030$  and  $0.577 \pm 0.027$  for Si and Ge, respectively, to be compared with our values of 0.548 and 0.548. The LDA [5] results of  $\zeta$  for Si and Ge are 0.56 and 0.57, respectively; they are in excellent agreement with the results of [17].

The ratios  $c_{12}/c_{11}$ ,  $[(c_{11} - c_{12})/2]/B \equiv \kappa$  and  $\beta/\alpha$  are also given in table 4. It is clear from the table that, according to our calculations, the  $\zeta$  and the ratio  $c_{12}/c_{11}$  monotonically

		Z*	β*		e <sub>T</sub> *		e <sub>p</sub> *
Material	вом	Experiment	вом	вом	Experiment	вом	Experiment
SiC	0.640	0.38	0.312	1.06	2.57	0.300	
BN	-0.096	0.55	0.429	0.476	2.47	-0.577	
BP	-0.452		0.269	0.093		-0.745	
BAs	-0.480		0.256	-0.139		-0.759	
AlP	0.350	0.30	0.598	1.15	2.28	-0.270	
AlAs	0.350	0.26	0.598	1.15	2.30	-0.270	-0.22
AlSb	0.157	0.19	0.530	0.864	1.93	-0.402	
GaP	0.279	0.24	0.574	1.04	2.04	-0.320	-0.28
GaAs	0.272	0.20	0.572	1.03	2.16	-0.325	-0.47
GaSb	0.051	0.15	0.498	0.703	2.15	-0.470	-0.42
InP	0.588	0.27	0.669	1.48	2.55	-0.086	
InAs	0.592	0.21	0.670	1.49	2.53	-0.083	-0.13
InSb	0.381	0.16	0.608	1.19	2.42	-0.248	-0.24
BeS	-0.451	_	0.658	0.427		-1.12	
BeSe	-0.452		0.658	0.426		-1.12	
BeTe	0.511		0.641	0.344		-1.17	
ZnS	0.198	0.42	0.767	1.22	2.15	-0.490	0.33
ZnSe	0.211	0.34	0.768	1.23	2.03	-0.475	0.13
ZnTe	0.168	0.28	0.766	1.19	2.00	-0.524	0.08
CdS	0.456		0.765	1.48	2.77	-0.176	
CdTe	0.442	0.32	0.766	1.46	2.35	-0.194	0.09
CuF	-0.805		0.767	0.218		-1.49	
CuCl	-0.445	0.30	0.756	0.563	1.12	-1.05	0.35
CaBr	-0.393	0.29_	0.750	0.607	1.49	-0.974	
Cuĭ	-0.294	0.21	0.734	0.684	2.40	-0.835	
AgI	-0.064	0.29	0.677	0.839	1.40	-0.495	

**Table 3.** Results calculated for diamond- and zincblende-structure semiconductors of effective atomic charge  $Z^*$ , transfer parameter  $\beta^*$ , transverse charge  $e_T^*$  and the piezoelectric charge  $e_p^*$ . For comparison, the values of experiment are also given.

decrease, while the ratios  $[(c_{11} - c_{12})/2]/B$  and  $\beta/\alpha$  monotonically increase with increase in covalency. The results versus covalency are shown in figures 1-4. The trend of  $\zeta$  as a function of  $\alpha_c$  is shown in figure 1; it is similar to that given by Martin [13], Harrison [1] and Kitamura and Harrison [10]. The trends of  $[(c_{11} - c_{12})/2]/B$  and  $c_{12}/c_{11}$  versus covalency  $\alpha_c$  are shown in figures 2 and 3, which are similar to that given by Kitamura *et al* [11]. Falter *et al* [16] have given the ratio of the non-central (bond-bending) force constant  $\beta$  to the central (bond-stretching) force constant  $\alpha$  as a function of polarity. This ratio is a measure of the importance of the covalent bond in stabilizing the tetrahedral structure. In figure 4, we give a plot of  $\beta/\alpha$  as a function of covalency  $\alpha_c$ . The trend is similar to that given by Falter *et al* [16].

(v) It is worth noting that there are several differences between the BOM and XHTB methods. One is the basis state chosen. In the BOM, sp<sup>3</sup> orbitals for cations and anions are selected as the basis states while, in the XHTB method, atomic orbitals are used as the basis states. A second is that the BOM method includes only the first-nearest-neighbour interaction for which universal parameters are optimized whereas, in the XHTB method, interactions are taken into account up to a sufficiently large distance that the interactions became negligibly small. The third is that the value of  $\lambda$  in the BOM method is a constant value of 0.738 for all the diamond- and zincblende-structure semiconductors, while its value in the XHTB method varies from 0.615 to 0.815. In addition, the bond length used in the BOM method is the theoretical value, while in the XHTB method it is the experimental value.

1-stretching force	so piven
Con)/2 and car bon	$V(B)$ and $R/\alpha$ are a
ess constants (c., -	K (≅ [(cu – cv)/2
r with elastic stiffne	the ratios c12/c11.
modulus B togethe	ment parameters ζ,
values of the bulk	's internal displace.
red by $(1 - \alpha_0^2)^{1/2}$ ,	istant $\beta$ . Kleinman
r covalency $\alpha_e$ defit	d-bending force cor
Table 4. Results fo	constant a and bon

	Øc	(10 <sup>11</sup> erg cm <sup>-3</sup> )	$(c_{11} - c_{12})/2$ (10 <sup>11</sup> erg cm <sup>-3</sup> )	$c_{44}$ (10 <sup>11</sup> erg cm <sup>-3</sup> )	α (10 <sup>3</sup> erg cm <sup>-2</sup> )	β (10 <sup>3</sup> ero cm <sup>-2</sup> )	*	cro lou	2	RIC
	001	60.1	205				~	111/211	×	<i>b/d</i>
) ;	0.1	1.50	0,00	40.9	230	62.6	0.548	0.410	0,486	0.264
5	1.00	5.01	5.08	7.10	52.3	13.8	0.548	0.410	0,486	0.264
.e	1.00	8,65	4.20	5.87	44.9	11.9	0.548	0.410	0.486	0.264
LS	1,00	4.40	2.14	2.99	26.2	6.92	0.548	0.410	0.486	0.264
SiC	0.987	25.9	12.5	17.5	107	28.1	0.551	0.414	0.481	0.262
BN	0.974	74.0	35.3	49.5	247	64.0	0.554	0.417	0.476	0.259
BP	166'0	25,9	12.5	17.5	107	28.2	0.550	0.413	0.482	0.262
BAs	0.992	22.4	10.8	15.1	95.7	25.1	0.550	0.413	0.483	0.263
AIP	0.942	9.31	4.31	6.11	46.5	11.6	0.563	0.427	0.463	0.251
AIAS	0.941	7.98	3.69	5.23	41.1	10.3	0.563	0.427	0,463	0.251
AISb	0.957	5.65	2.65	3.74	31,4	7.99	0,558	0.422	0,470	0.255
GaP -	0.947	8.94	4,14	5.89	45.2	11.3	0.561	0.425	0.466	0.251
GaAs	0.948	7.69	3.58	5.07	40.0	10.1	0.561	0.425	0.466	0.252
GaSb	0,965	5.49	2.59	3.65	30.8	7.88	0.556	0.420	0.474	0.256
	0.918	5.89	2.67	3,80	31.9	7.82	0.570	0.435	0.453	0.245
InAs	0.917	5.04	2.28	3.25	28.2	6.89	0.570	0.435	0.453	0.245
usb 1	0.937	3.66	1.69	2.40	22.0	5.50	0.563	0.428	0.462	0.250
BeS	0.922	25.3	11.5	16.4	102	25,2	0.568	0.434	0,455	0,246
BeSe	0.922	20.2	9.18	13.1	85.6	21.1	0.568	0,434	0,455	0.246
Beffe	0.928	12.6	5.77	8.21	59.0	14.6	0.566	0.432	0,457	0.247
SUS 6	0.836	7.78	3.20	4.69	38.5	8.51	0.598	0.468	0.412	0.221
Zuse	0.833	6.12	2.51	3.68	31,8	6.99	0.599	0.469	0.411	0.220
Zule	0.840	3.84	1.59	2.33	21.9	4.88	0.596	0.466	0.415	0.222
	0,789	4.79	1.84	2.75	25.7	5.26	0.618	0.491	0.385	0.205
Cdie	0.792	2.33	0.90	1.34	14.4	2.98	0.616	0.490	0.387	0.207
ы Сбғ	0.836	59.3	24,4	35.8	195	43.3	0.598	0.468	0,412	0.221
ה בפכו	0,769	9.87	3,68	5.53	45.4	9.02	0.627	0.502	0.372	0.199
CuBr C	8 <i>2</i> 7.0	6.88	2.51	3.79	33.8	6.57	0.632	0.509	0.365	0.194
5	0.736	3.44	1.21	1.84	19.3	3.60	0.644	0.522	0.350	0.187
Agi	0.679	1.96	0.61	0.95	12.0	1.97	0,677	0.563	0.308	0,164

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Figure 1. Kleinman's internal displacement parameters  $\zeta$  versus covalency  $\alpha_c$ .

Figure 2. Plot of  $\kappa \equiv [(c_{11} - c_{12})/2]/B)$  versus covalency  $\alpha_c$  listed in table 4.



Figure 3. Curve for the ratio  $c_{12}/c_{11}$  as a function of covalency  $\alpha_c$ .



Figure 4. Ratio  $\beta/\alpha$  of bond-bending constant  $\beta$  to central-force constant  $\alpha$  as a function of the covalency  $\alpha_c$ .

Comparison of bulk modulus B and elastic stiffness constants  $(c_{11} - c_{12})/2$ , and  $c_{44}$ , obtained from BOM calculations with the experimental values and XHTB calculation results [11] are shown in table 5. The bond-stretching force constant  $\alpha$  and bond-bending force constant  $\beta$  obtained from BOM calculations (neglecting the Coulomb contribution) with the theoretical quantities  $\alpha$  and  $\beta$  derived from the experimental values of B and  $(c_{11} - c_{12})/2$  (the Coulomb contribution is included) are also given in the table. Table 5 indicates that the results obtained by using the BOM are in good agreement with the experimental and other theoretical results.

(vi) Table 6 shows the LDA results together with the experiment and our calculations for bulk Si, Ge, AlP, AlAs, GaP and GaAs. Comparing the LDA results with experiment, we can see that the LDA results agree very well. The excellent agreement between LDA results and experimental data prove that the LDA method is accurate. Although the BOM results are not as accurate as the LDA, the BOM results are reasonable. Because the BOM is very

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				mono	outer results, Life une	יטגכוורמו ו	huammus	a alle p wele uch	rea Irolli	nie exper	intental values (	or <i>b</i> and	(c11 — c12)/2.
		B (×10 <sup>11</sup> erg cm <sup>-3</sup> )		- 113)	c12)/2 (×10 <sup>11</sup> erg (	cm <sup>-3</sup> )	ŭ	44 (x10 <sup>11</sup> erg cm <sup>-3</sup>		α (×1(	$3^3 \text{ erg cm}^{-2}$ )	β (×1	) <sup>3</sup> erg cm <sup>-2</sup> )
	BOM	Experiment [18]	ХНТВ	BOM	Experiment [18]	XHTB	BOM	Experiment [18]	XHTB	BOM	Theory [13]	BOM	Theory [13]
J	69.1	56.7	39.3	33.6	26.5	35.0	46.9	43.0	35.4	236		62.6	
		44.2			47.6			57.7			129.33		84.76
Si	10.46	9.78	8.70	5.08	5.09	7.14	7.10	7,96	7.83	52.3	48.50	13.8	13.81
Ge Ge	8.65	7.52	7.63	4.20	4.03	5.49	5,87	6.71	6.87	44.9	38.67	11.9	11.35
Sn	4.40	5.31	4.60	2.14		3.43	2,99		4.14	26.2	25.45 [19]	6.92	6.44 [19]
SiC	25.9	22.8	18.5	12.5	10.5	9.60	17.5	14.9	13.6	107	88.00 [19]	28.1	47.52 [19]
BPe	25.9	17.2	16.5	12.5	10.8	14.2	17.5	16.0	14.4	107	,	28.2	
AIP	9.31	8.60	8.05	4.31		3.55	6,11		5.20	46.5	47,29 [19]	11.7	9.08 [19]
ALAs	7.98	7.73	7.40	3.69		3.77	5.23		5.14	41.1	43.05 [19]	10.3	9.86 [19]
AISb	5.65	5.93	5.41	2,65	2.26	3,36	3,74	4,16	4.30	31.4	35.35	7.99	6.77
GaP	8.94	8.87	8.15	4.14	3.94	3.55	5,89	7.05	5.56	45.2	47.32	11.3	10.44
GaAs	7.69	7.48	7.24	3.58	3.25	3.25	5.07	5.92	5.29	40.0	41,19	10.1	8.95
GaSb	5.49	5.63	5.54	2.59	2.41	3.28	3.65	4.32	4.52	30.8	33.16	7.88	7.22
ЧП	5.89	7.25	6.08	2.67	2.23	2.01	3,80	4,60	3.75	31.9	43.04	7.82	6.24
InAs	5,04	5.80	5.63	2.28	1.90	2.16	3,25	3.96	3.73	28.2	35,15	6.89	5.50
lnSb	3.66	4.66	4.40	1.69	1.51	2.14	2,40	3.02	3.35	22.0	29.61	5.50	4.77
ZnS	7.78	7.80	7.33	3,20	1.95	1.41	4,69	4.62	2.49	38.5	44.92	8.51	4.78
ZnSe	6.12	5.95	6.51	2.51	1.61	1.56	3,68	4,41	2.64	31.8	35,24	6.99	4.23
ZnTe	3.84	5.09	5,19	1.59	1.53	1.80	2.33	3.12	2.75	21.9	31.35	4.88	4.45
CdS	4.79	6.44	5.32	1.84		0.71	2.75		1.64	25.7		5.26	
CdTe	2.33	4.24	4.02	06.0	0.84	1.12	1.34	1.99	1.92	14.4	29.02	2.98	2,43

Table 6	. Compari	ison of bulk moduli J	B and elastic	stiffness col	nstants (c	c11 - c12)/2 and c44	obtained fro	om LDA calci	ulations a	ud experimental and	BOM calcul	ational results.
		B (×10 <sup>11</sup> erg	; cm <sup>-3</sup> )			$(c_{11} - c_{12})/2$ (×10	11 erg cm <sup>-2</sup>	Ĵ.		c44 (×10 <sup>11</sup> er	g cm <sup>-3</sup> )	
	BOM	Experiment [18]	LDA [5]	[/] Vari	BOM	Experiment [18]	LDA [5]	[7] AGJ	BOM	Experiment [18]	(2) Van	[ <u>]</u> ] VOT
Si	10.46	9.78	9.70	9.30	5.08	5.09	5.00	4.90	7,10	7.96	8.00	8.50
පී	8.65	7.52	7.70	7.20	4.20	4.03	4,25	4.25	5.87	6.71	6,90	6.30
AIP	9.31	8.60	8.65 [6]		4.31				6.11			
AIAs	7.98	7.73	7,41 [6]		3.69				5.23			
GaP	8.94	8.87	8.97 [6]		4.14	3.94			5.89	7.05		
GaAs	7.69	7.48	7.25 [6]	7.30	3.58	3.25		3.50	5.07	5.92		6.20

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simple, it can be used for the systematic study of structural properties in many materials and obtains the structural trends.

# 4. Summary

We have calculated the bond length d, bond polarity  $\alpha_p$ , bulk modulus B, elastic shear constant  $(c_{11}-c_{12})/2$  and  $c_{44}$ , bond-stretching force constant  $\alpha$ , bond-bending force constant  $\beta$ , internal displacement parameter  $\zeta$ , effective atomic charge  $Z^*$ , transfer parameter  $\beta^*$ , transverse charge  $e_T^*$  and piezoelectric charge  $e_p^*$  for 30 diamond- and zincblende-structure semiconductors using the BOM and compared these results with values obtained from other theoretical calculations and experiments. For most materials, we find the following.

(1) The bond length and polarity are in good agreement with the experiments.

(2) The transverse charges are about one electron less than the experiments, and the piezoelectric charges are in good agreement with the experiments (for III-V compound semiconductors).

(3) The values of bulk modulus and elastic stiffness constants obtained from BOM calculations are in good agreement with the experimental and the XHTB and LDA theoretical results.

(4)  $\zeta$  and the ratio  $c_{12}/c_{11}$  monotonically decrease, while the ratios  $[(c_{11} - c_{12})/2]/B$  and  $\beta/\alpha$  monotonically increase with increasing covalency. Their trends are the same as experimental and other theoretical results.

# References

- [1] Harrison W A 1980 Electronic Structure and the Properties of Solids (New York: Freeman)
- [2] Baranowski J M 1984 J. Phys. C: Solid State Phys. 17 6287
- [3] Chadi D J and Martin R M 1976 Solid State Commun. 19 643
- [4] Harmon B N, Weber W and Hamann D R 1982 Phys. Rev. B 25 1109
- [5] Wei Siqing, Allan D C and Wilkins J W 1992 Phys. Rev. B 46 12411
- [6] Froyen S and Cohen M L 1983 Phys. Rev. B 28 3258
- [7] Nielsen O H and Martin R M 1985 Phys. Rev. B 32 3792
- [8] Christensen N E 1984 Solid State Commun. 50 177
- [9] Methfessel M, Rodriguez C O and Ahdersen O K 1989 Phys. Rev. B 40 2009
- [10] Kitamura M and Harrison W A 1991 Phys. Rev. B 44 7941
- [11] Kitamura M, Muramatsu S and Harrison W A 1992 Phys. Rev. B 46 1351
- [12] Herman F and Skillman S 1963 Atomic Structure Calculations (Englewood Cliffs, NJ: Prentice Hall)
- [13] Martin R M 1970 Phys. Rev. B 1 4005
- [14] Mitra S S and Massa N E 1982 Handbook on Semiconductors vol 1, ed W Paul (Amsterdam: North-Holland) p 81
- [15] Berding M A, Sher A and Chen A-B 1987 Phys. Rev. B 36 7433
- [16] Falter C, Ludwig W, Selmke M and Zierau W 1984 Phys. Lett. 105A 139
- [17] Molinas-Mata P and Cardona M 1991 Phys. Rev. B 43 9799
- [18] Muramatsu S and Kitamura M 1993 J. Appl. Phys. 73 4270
- [19] Martins J L and Zunger A 1984 Phys. Rev. B 30 6271